

Scordelis-Lo Roof Revisited

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Shell Finite Elements

 The error estimates for standard finite elements for shells are of the form

$$\frac{||\mathbf{u} - \mathbf{u}_{h,p}|||}{|||\mathbf{u}|||} \le C \frac{L}{t} \left(\frac{h}{L}\right)^p$$

• There are formulations (like MITC-type) for which

$$\frac{|||\mathbf{u} - \mathbf{u}_{h,p}|||_{h}}{|||\mathbf{u}|||} \le C\left(\frac{h}{L}\right)^{p}$$

under very specific geometric assumptions.

Practical shell finite element formulations are validated using benchmark computations



Scordelis-Lo Roof Problem Setup



Problem Setup (SI-units)



 $R = 25 \,\mathrm{m}$ $L = 50 \,\mathrm{m}$ $\alpha = 40^{\circ}$



Problem Setup (cont.) (SI-units)

- Material:
 - Isotropic material with Young's modulus set to $E = 4.32 \cdot 10^8 \text{ N/m^2}$.
 - Poisson's ratio set to $\nu = 0.0$.
- Kinematic constraints:
 - The outer curved edges are constrained against vertical translations (*diaphragm support*).
 - The outer straight edges are *free*.
- Load:
 - A force per are unit of −90 N/m² in the vertical direction is applied on the surface.



Benchmark Setting

- Set the thickness to $t = 0.25 \,\mathrm{m}$ and determine the value of the vertical displacement at the center of the free edge.
- The reference solution quoted in the literature for the midside vertical displacement is -0.309 m.

Scordelis, A. C., & Lo, K. S. (1964). Computer Analysis of Cylindrical Shells. *ACI Journal Proceedings*, *61*(5). doi:10.14359/7796



Classical Shell Model (Love-Novozhilov-Koiter)

Total potential energy of the shell is

$$\mathcal{F}(\boldsymbol{u}) = \frac{Et}{2} \int_{\omega} (\beta_{11}^2 + 2\beta_{12}^2 + \beta_{22}^2) \, \mathrm{d}x \mathrm{d}\phi + \frac{Et^3}{24} \int_{\omega} (\kappa_{11}^2 + 2\kappa_{12}^2 + \kappa_{22}^2) \, \mathrm{d}x \mathrm{d}\phi$$
$$- p \int_{\omega} (\sin \phi \cdot v - \cos \phi \cdot w) \, \mathrm{d}x \mathrm{d}\phi,$$

where $\boldsymbol{u} = (u, v, w)$ and $\omega = (0, L) \times (-\alpha, \alpha)$.

• The membrane strains are defined as

$$\beta_{11} = \frac{\partial u}{\partial x}, \quad \beta_{22} = \frac{1}{R} \left(\frac{\partial v}{\partial \phi} + w \right), \quad \beta_{12} = \frac{1}{2} \left(\frac{1}{R} \frac{\partial u}{\partial \phi} + \frac{\partial v}{\partial x} \right)$$



Classical Shell Model (cont.) Bending Strains

Shallow shell theory (LNK-SS):

$$\kappa_{11} = \frac{\partial^2 w}{\partial x^2}, \quad \kappa_{22} = \frac{1}{R^2} \frac{\partial w}{\partial \phi^2}, \quad \kappa_{12} = \frac{1}{R} \frac{\partial^2 w}{\partial x \partial \phi}.$$

• Novozhilov's model (LNK-S):

$$\kappa_{11} = \frac{\partial^2 w}{\partial x^2}, \quad \kappa_{22} = \frac{1}{R^2} \frac{\partial w}{\partial \phi^2} - \frac{1}{R^2} \frac{\partial v}{\partial \phi}, \quad \kappa_{12} = \frac{1}{R} \frac{\partial^2 w}{\partial x \partial \phi} - \frac{1}{R} \frac{\partial v}{\partial x}.$$

• Vlasov's model (Vlasov):

$$\kappa_{11} = \frac{\partial^2 w}{\partial x^2}, \quad \kappa_{22} = \frac{1}{R^2} \frac{\partial w}{\partial \phi^2} + \frac{w}{R^2}, \quad \kappa_{12} = \frac{1}{R} \frac{\partial^2 w}{\partial x \partial \phi} - \frac{1}{2R} \frac{\partial v}{\partial x} + \frac{1}{2R^2} \frac{\partial u}{\partial \phi}.$$



Semi-analytic Solution



Ansatz

• Kinematic constraints at x = 0, L are satisfied for

$$u = \sum_{m=1}^{\infty} U_m(\phi) \cos(\beta_m x), \ v = \sum_{m=1}^{\infty} V_m(\phi) \sin(\beta_m x), \ w = \sum_{m=1}^{\infty} W_m(\phi) \sin(\beta_m x),$$

where

$$\beta_m = \frac{m\pi}{L}$$

• The constant load can be represented as

$$p = \sum_{m=1}^{\infty} P_m \sin(\beta_m x), \quad P_m = \begin{cases} \frac{4p}{\pi m}, & m \text{ odd} \\ 0, & m \text{ even} \end{cases}$$



Euler Equations Shallow Shell Theory

• Substitution of the Ansatz in $\mathcal{F}(u)$ and minimization the energy with respect to U_m, V_m, W_m gives

$$\beta_m^2 U_m - \frac{\beta_m}{2R} V'_m - \frac{U''_m}{2R^2} = 0,$$
$$\frac{\beta_m^2}{2} V_m + \frac{\beta_m}{2R} U'_m - \frac{1}{R^2} (V''_m - W'_m) = \frac{P_m}{Et} \sin \phi,$$
$$\frac{1}{R^2} (V'_m + W) + \frac{t^2}{12} \left(\frac{W_m^{(4)}}{R^4} - \frac{\beta_m^2}{2R^2} W''_m + \beta_m^4 W_m \right) = -\frac{P_m}{Et} \cos \phi.$$

• These constitute a system of *constant-coefficient* ODEs of total degree *eight,* valid when $\phi \in (-\alpha, \alpha)$.

Natural Boundary Conditions Shallow Shell Theory

• The natural boundary conditions at the free edges $a = \pm \alpha$ are

$$\frac{1}{R}U'_{m} + \beta_{m}V_{m} = 0,$$

$$V'_{m} + W_{m} = 0,$$

$$\beta_{m}^{2}W'_{m} - \frac{1}{2R^{2}}W_{m}^{(3)} = 0,$$

$$W''_{m} = 0.$$

 These correspond to the conditions of vanishing normal forces, effective shear force and bending moment.



Solution

• The solution of the system can be expressed as $\frac{1}{8}$

$$(U_m, V_m, W_m) = (\hat{U}_m, \hat{U}_m, \hat{U}_m) + \sum_{i=1}^{\circ} C_i \boldsymbol{v}_i e^{\lambda_i \phi},$$

where

- $(\hat{U}_m, \hat{U}_m, \hat{U}_m)$ is a special solution of the *non-homogeneous* system.
- (λ_i, v_i) are the eigenpairs associated to the *homogeneous* system.
- $(\hat{U}_m, \hat{U}_m, \hat{U}_m)$ is found easily in the form $(\hat{U}_m(\phi), \hat{U}_m(\phi), \hat{U}_m(\phi)) = (\tilde{U}_m \cos \phi, \tilde{V}_m \sin \phi, \tilde{W}_m \cos \phi)$



Solution (cont.)

- The special solution and the eigenvalues can be determined by symbolic computation for each m.
- The eigenvectors are solved numerically.
- The most difficult part is the determination of the integration coefficients $C_i, i = 1, ..., 8$.
- These can be solved with Mathematica using arbitrary high numerical precision.



Numerical Results



Midpoint Vertical Displacement

• Original problem setup with t=0.25 m (R/t=100):

LNK-SS	LNK-S	Vlasov	COMSOL	MITC4S
-0.309 m	-0.299 m	-0.312 m	-0.302 m	-0.302 m

Academic problem setup with t=0.025 m (R/t=1000):

LNK-SS	LNK-S	Vlasov	COMSOL	MITC4S
-30.9 m	-32.0 m	-32.0 m	-32.0 m	-32.0 m



In-plane Shear Force LNK-S vs. COMSOL (4-node, normal)





In-plane Shear Force LNK-S vs. COMSOL (4-node, extra-fine)





Concluding Remarks

- We have presented a semi-analytic solution technique for the classical Scordelis-Lo roof model problem.
- Our approach allows to study different variants of shell models for different values of the thickness.
- The solution technique can be used to benchmark shell finite elements but also directly in the design of cylindrical shell roofs!

