Global analysis of active longitudes of sunspots

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ABSTRACT

Context. Active longitudes have been found in various manifestations of solar activity. The longitudinal distribution of, e.g., sunspots and solar X-ray flares shows two persistent preferred longitudes separated by roughly 180 degrees. We previously studied solar X-ray flares using an improved version of a dynamic, differentially rotating coordinate system and found enhanced rotational asymmetry and rotation parameter values that are consistent for the three classes of X-ray flares.

Aims. We aim to find the optimal values of rotation parameters of active longitudes of sunspots for several different time intervals and separately for the two solar hemispheres.

Methods. We perform a global study of the longitudinal location of sunspots (all sunspots and first appearance sunspots) using a refined version of a dynamic, differentially rotating coordinate system.

Results. We find that the rotation parameters for sunspots are in good agreement with those obtained for X-ray flares using the same method. The improved method typically finds somewhat faster equatorial rotation with better accuracy. The improved treatment also leads to a larger non-axisymmetry. Rotation parameters for all sunspots and first appearances closely agree with each other, but non-axisymmetry is systematically larger for all sunspots than for first appearances, suggesting that strong fields follow more closely the pattern of active longitudes. The refined method emphasizes hemispheric differences in rotation. Over the whole interval, the mean rotation in the southern hemisphere is slower than in the north. We also find significant temporal variability in the two rotation parameters over the 136-year interval. Interestingly, the long-term variations (trends and residual oscillations) in solar rotation are roughly the opposite in the northern and southern hemispheres.

Conclusions. Rotation parameters vary differently with time in the northern and southern hemispheres. Both sunspots and flares strongly suggest that the northern hemisphere rotated considerably faster but the southern hemisphere slightly slower than the Carrington rotation rate during the last three solar cycles.

Key words. Sun: rotation - sunspots

1. Introduction

Active longitudes of the Sun have been studied in detail since the early 20th century (Chidambara 1932; Losh 1939). It has been found that various manifestations of solar activity, such as sunspots (Berdyugina et al. 2003; Usoskin et al. 2005; Juckett 2006), solar magnetic field (Benevolenskaya et al. 1999; Bumba et al. 2000), heliospheric magnetic field (Ruzmaikin et al. 2001; Takalo & Mursula 2002; Mursula & Hiltula 2004), and flares (Bai 1988, 2003) occur preferentially at specific longitudes that are called active longitudes. However, lifetimes and rotation velocities of active longitudes obtained for different tracers of solar activity have been found to vary significantly. Balthasar (2007) noted, by analyzing sunspot numbers (1818-2006), that the rotation period varies over long timescales. Heristchi & Mouradian (2009) found an increase in the rotation rate from cycle 19 to cycle 23 by calculating sunspot numbers and the coronal radio flux at 10.7 cm. Pelt et al. (2010), using a nonparametric statistical analysis of the longitude of sunspot groups in 1874–2008, proposed that the strong tendency for sunspot groups to cluster at a certain longitude peters out with time, the longest observable correlations reaching 10-15 Carrington rotations. This result, however, is more relevant for active regions than for the long-living active longitudes studied here.

Using a dynamic, differentially rotating coordinate system, Usoskin et al. (2005) found that two active longitudes of sunspots separated by about 180° existed for the whole time interval of sunspots with measured locations (1878–1996). The same system was used to analyze the longitudinal occurrence of solar X-ray flares observed by the NOAA/GOES satellites (Zhang et al. 2007a,b). However, in these studies the differential rotation parameters for sunspots were quite different from those obtained for X-ray flares. Even the parameters for the different X-ray class flares when calculated separately were found to be quite different (Zhang et al. 2007b).

We have recently conducted a global search of the optimized rotation parameters for the three classes of solar X-ray flares using a corrected version of the dynamic, differentially rotating coordinate system (Zhang et al. 2011). We found that the refined parameters describing the rotation of active longitudes are mutually consistent for the three flare classes and that the northern hemisphere rotates significantly faster than the southern hemisphere during the last three solar cycles. This contradicts the result of Heristchi & Mouradian (2009), who found no significant difference between the rotation rates of the two hemispheres. Moreover, the refined parameters yielded a consistently higher level of non-axisymmetry for the longitudinal distribution, thus increasing evidence for the existence of active longitudes. Here we perform a global search to find the refined optimal values of parameters describing the longitudinal distribution of sunspots, all sunspots and the first appearance of sunspots separately, using the corrected version of the dynamic, differentially rotating coordinate system. We obtain the global best-fit values of parameters describing the differential rotation of active longitudes of sunspots, and estimate their statistical errors and the level of non-axisymmetry. We compare these values for sunspots with the corresponding values of parameters found for solar X-ray flares.

The paper is organized as follows. The data is introduced in Sect. 2. The analysis method based on the corrected version of the dynamic, differentially rotating coordinate system is described in Sect. 3. The results obtained for the differential rotation parameters and non-axisymmetry are presented in Sects. 4 and 5. The results are discussed in Sect. 6. In Sect. 7 we present the final conclusions of this study.

2. Data

We use here the sunspot group data collected by the Royal Greenwich Observatory and the USAF/NOAA for the years 1874-2009 (http://solarscience.msfc.nasa.gov/ greenwch.shtml). We use this data of sunspot group locations and areas in two different ways. Firstly, only the first appearance of each sunspot group is taken into account, i.e., each sunspot group is considered only once when it is first mentioned in the data (either on the day of its birth or when it first appeared in the East limb), ignoring all later records of the spot. It is important to study the newly appeared sunspots because they are considered to better indicate the emergence of magnetic flux through the photosphere, which is a key process of the solar cycle (Solanki 2003) and of the study of the solar magnetic dynamo. Secondly, all daily records of sunspots are taken into account, irrespective of whether they are new or old. We also note that we use here data that are uncorrected for the change in data source in 1977 and ignore the area correction factor of 1.4 since 1977 (see above web page). However, this change in the homogeneity of the database has little importance for the type of studies conducted here.

3. Improved analysis method

Differential rotation of the solar surface can be described approximately as

$$\Omega_{\phi} = \Omega_0 - B\sin^2\phi, \tag{1}$$

where Ω_{ϕ} (deg/day) is the sidereal angular velocity at a given latitude ϕ , Ω_0 (deg/day) is the (sidereal) equatorial angular velocity and *B* (deg/day) describes the differential rotation rate. Rotation parameters Ω_0 and *B* are constants whose values will be determined by optimization.

Taking the Carrington rotation period as the time step, the normalized area of a *k*th sunspot group in the *j*th rotation is

$$W_{jk} = A_{jk} / \sum_{k} A_{jk}, \tag{2}$$

where A_{jk} is the observed sunspot area (corrected for projection), and the sum is taken over all spots in the *j*th rotation. The sunspot area-weighted latitude for this rotation is then

$$\langle \phi_j \rangle = \sum_k W_{jk} \phi_{jk},\tag{3}$$

where ϕ_{jk} denotes the observed latitude of the *k*th spot in the *j*th rotation. For Carrington rotations with no sunspots, we use the linear interpolation of the mean latitude. Accordingly, the (sidereal) angular velocity for the *j*th rotation is

$$\Omega_{N_j} = \Omega_0 - B \sin^2 \langle \phi_j \rangle. \tag{4}$$

On the *k*th day of the *i*th Carrington rotation, the migration M of an active longitude with respect to the Carrington reference frame is determined by

$$M = T_{\rm c} \sum_{j=N_0}^{N_{i-1}} (\Omega_j - \Omega_{\rm c}) + k(\Omega_{N_i} - \Omega_{\rm c}),$$
(5)

where N_0 and N_{i-1} denote the Carrington rotation numbers of the first and the (i-1)th rotation of the data set, $T_c = 27.2753$ days is the synodic Carrington rotation period, Ω_c is the angular velocity of Carrington frame (in sidereal frame 14.1844 deg/day and in synodic frame 13.199 deg/day), and *k* is the time of sunspot group observation given as a fractional day of the Carrington rotation.

Let us assume that one active longitude in the beginning is at Carrington longitude Λ_{01} . It will reach Λ_{ik1} on the *k*th day of the *i*th Carrington rotation, which can be expressed as

$$\Lambda_{ik1} = (\Lambda_{01} + M) \mod 360^\circ, \tag{6}$$

where mod 360° means modulo 360° within the range [0°, 360°]. The other active longitude is assumed to be at $\Lambda_{ik2} = \Lambda_{ik1} \pm 180^\circ$, i.e., opposite to Λ_{ik1} .

The deviation Δ_{ik} between the longitudinal position λ_{ik} of a sunspot, which occurred on the *k*th day in the ith rotation, and one of the two active longitude bands (Λ_{ik1} or Λ_{ik2}) is

$$\Delta_{ik} = \min(|\lambda_{ik} - \Lambda_{ik1}|, \ 360^{\circ} - |\lambda_{ik} - \Lambda_{ik1}|, \ |180^{\circ} - |\lambda_{ik} - \Lambda_{ik1}||).$$
(7)

In order to illustrate this, we discuss some examples. Assuming $\Lambda_{ik1} = 2^{\circ}$ and $\Lambda_{ik2} = 182^{\circ}$, and taking $\lambda_{ik} = 5^{\circ}$, one obtains $\Delta_{ik} = |\lambda_{ik} - \Lambda_{ik1}| = 3^{\circ}$. On the other hand, taking $\lambda_{ik} = 359^{\circ}$, yields correspondingly $\Delta_{ik} = 360^{\circ} - |\lambda_{ik} - \Lambda_{ik1}| = 3^{\circ}$. When $\lambda_{ik} = 179^{\circ}$ or $\lambda_{ik} = 185^{\circ}$, $\Delta_{ik} = |180^{\circ} - |\lambda_{ik} - \Lambda_{ik1}| = 3^{\circ}|$.

We note that the first studies using the dynamic reference frame (Usoskin et al. 2005; Zhang et al. 2007b) missed the second term on the right-hand side of Eq. (5). Thus, the mean migration with respect to the Carrington rotation (the first term on the right-hand side of Eq. (5)) was used as the migration of active longitudes from the beginning to the end of the rotation, which resulted in only one active longitude Λ_{i1} in the *i*th rotation being used for each spot of that rotation when defining the deviation for spots according to Eq. (7). Adding the second term of Eq. (5), we obtain a more accurate, smoothly changing active longitude Λ_{ik1} at the time when a spot occurs, and a more precise deviation Δ_{ik} in Eq. (7) by calculating $|\lambda_{ik} - \Lambda_{ik1}|$, rather than $|\lambda_{ik} - \Lambda_{i1}|$. This improved recipe was first adopted in a recent study of the active longitudes of solar X-ray flares (Zhang et al. 2011).

We define the merit function (which is to be minimized) as the weighted mean squared deviation of sunspots from the nearest active longitude

$$\epsilon = \frac{1}{n} \sum_{i} \sum_{k} W_{jk} \Delta_{ik}^{2}, \tag{8}$$

where n denotes the total number of sunspots and W_{jk} is defined in Eq. (2). We first perform a global search of the best-fit parameters as the global minimum of the merit function by

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Fig. 1. Values of the merit function in color coding (in units of deg²) as a function of Ω_0 and *B* for all sunspots of the southern hemisphere in the last three solar cycles (**a**, *top*) over the full studied parameter range; (**b**, *bottom*) in the limited region around the best-fit parameter values.

varying the value of Ω_0 in the interval [13.5, 15.0] (deg/day) in steps of 0.01 (deg/day), *B* within [0.0, 5.0] (deg/day) in steps of 0.01 (deg/day) and Λ_{01} within [0°, 360°] in steps of 1°. After finding the global minimum, we perform another round of optimization over a more limited parameter range using steps of 0.001 (deg/day) for both Ω_0 and *B*. We always analyze the sunspots in the northern and southern hemispheres separately. We also define all parameters for two data sets, one using all sunspot groups included in the above mentioned database and the other using the first appearance of spots only. We also use different time intervals, varying the length of intervals from one solar cycle to three cycles and six cycles, as well as using the 129-year interval (1878–1996) of Usoskin et al. (2005) and the whole 136-year period of data in 1874–2009.

Figure 1 shows the values of the merit function in color coding as a function of Ω_0 and *B* for all sunspots in the southern hemisphere during the last three solar cycles. Figure 1a depicts the global situation over the whole parameter range, showing the global minimum of the merit function as the only region where the lowest values (depicted by dark blue color) are found. The striped structure of the merit function in Fig. 1a (and all similar plots) is a result of the correlation of the two parameters and the different sensitivity of the merit function to them. Figure 1b shows the detailed situation of the merit function in a more limited range (the area corresponding to the smallest bin value of merit function in Fig. 1a) around the best-fit parameters found in the global search, using the smaller step size of 0.001 (deg/day) for Ω_0 and *B*. Finding the minimum of the merit function in the reduced range by using a finer step size we can calculate the best-fit parameters more accurately. Moreover, detailed plots such as Fig. 1b were used to obtain the errors in the two parameters. As depicted in Fig. 1b, the values of the merit function were divided to eight equal bins between the largest and the smallest value within the range included in Fig. 1b. The parameter errors were then calculated as corresponding to the standard deviations of the merit function in either direction within the region of its smallest value (denoted in dark blue in Fig. 1b).

4. Differential rotation parameters

We have listed the optimal values of the differential rotation parameters with error bars for the different time intervals in Tables 1–4. In addition to the parameters Ω_0 , these tables also include the values of Ω_{17} , the rotation angular velocity at the latitude of 17°, which is the mean latitude of sunspots. In each table we have treated the two hemispheres separately and calculated the optimal parameters both using all sunspots and the first appearances of spots. Table 1 gives the parameter values and their errors for each solar cycle separately. Table 2 lists the same parameters for three-cycle time intervals and Table 3 for six-cycle intervals. Finally, Table 4 depicts the parameters for the whole time interval 1974–2009, as well as for 1874–1996, allowing comparison with the earlier results (Usoskin et al. 2005).

Figure 2 shows the merit function for a detailed analysis of all sunspots observed in the southern hemisphere during solar cycle 21 (Fig. 2a), during the last six solar cycles (Fig. 2b) and for the whole 136-year interval (Fig. 2c). Obviously, both the value and the accuracy of the fit parameters depend on the length of the analyzed period, with less accurate values obtained for shorter lengths. For solar cycle 21 (Fig. 2a), the optimal equatorial angular velocity (see Table 1) is $\Omega_0 = 14.240 \pm$ 0.018 (deg/day), while for the last six solar cycles, solar cycles 18–23 (see Table 3) it is $\Omega_0 = 14.322 \pm 0.011$ (deg/day) and for the whole interval $\Omega_0 = 14.401 \pm 0.004$ (deg/day). The improved accuracy in the two parameters is seen as the narrowing and shortening of the best-fit region in Fig. 2. However, we note that the accuracy does not improve very quickly when increasing the time interval from one cycle to six cycles or 136 years because the different cycles rotate slightly differently. The average error bars of the best-fit parameters Ω_0 , *B*, and Ω_{17} denoted by E_{Ω_0}, E_B , and $E_{\Omega_{17}}$ are given in Table 5 for the different time intervals. The average errors decrease slowly but systematically with increasing time interval. For instance, the mean error bars of Ω_{17} for the first appearance of spots decrease from ± 0.005 (deg/day) for 1-cycle values to ± 0.001 (deg/day) for 136-year values in both hemispheres. The same decrease is seen for all parameters and both hemispheres.

In Fig. 3 we have depicted the values of Ω_0 using all sunspots for the individual cycles and three-cycle intervals (from Tables 1 and 2). The two upper (lower) panels show the results for the northern (southern) hemisphere, while the left panels depict the individual cycles, and the right panels the three-cycle periods. It is interesting to see how much the optimal values of Ω_0 vary among the individual cycles and even between the two hemispheres of one cycle. The fastest equatorial rotation is seen during cycle 14 and the slowest in cycle 22, both in the southern hemisphere. The slowest and fastest equatorial rotations in the northern hemisphere are seen during the successive cycles 20

	Northern hemisphere								
	Firs	t appearance of s		All spots					
cycle	Ω_0	В	Ω_{17}	Ω_0	В	Ω_{17}			
12	14.280 ± 0.030	1.547 ± 0.372	14.151 ± 0.006	14.283 ± 0.010	1.505 ± 0.133	14.153 ± 0.003			
13	14.209 ± 0.017	3.497 ± 0.241	13.910 ± 0.005	14.200 ± 0.016	3.319 ± 0.243	13.915 ± 0.005			
14	14.226 ± 0.021	1.166 ± 0.245	14.125 ± 0.004	14.242 ± 0.016	1.664 ± 0.241	14.100 ± 0.005			
15	14.387 ± 0.011	3.828 ± 0.153	14.060 ± 0.003	14.346 ± 0.012	3.021 ± 0.172	14.086 ± 0.003			
16	14.429 ± 0.017	3.726 ± 0.204	14.110 ± 0.004	14.433 ± 0.020	3.827 ± 0.236	14.106 ± 0.004			
17	14.431 ± 0.033	4.223 ± 0.412	14.068 ± 0.003	14.437 ± 0.007	4.033 ± 0.077	14.092 ± 0.001			
18	14.369 ± 0.022	0.406 ± 0.260	14.331 ± 0.004	14.374 ± 0.015	0.467 ± 0.163	14.331 ± 0.003			
19	14.309 ± 0.016	1.524 ± 0.156	14.179 ± 0.004	14.292 ± 0.016	1.347 ± 0.158	14.177 ± 0.003			
20	14.203 ± 0.015	2.515 ± 0.200	13.988 ± 0.003	14.188 ± 0.009	2.340 ± 0.115	13.988 ± 0.002			
21	14.595 ± 0.030	3.435 ± 0.357	14.302 ± 0.004	14.599 ± 0.025	3.780 ± 0.289	14.273 ± 0.003			
22	14.495 ± 0.010	2.162 ± 0.098	14.309 ± 0.002	14.497 ± 0.024	2.090 ± 0.254	14.318 ± 0.005			
23	14.511 ± 0.031	0.421 ± 0.481	14.467 ± 0.010	14.530 ± 0.034	0.659 ± 0.448	14.471 ± 0.008			
	Southern hemisphere								
		t appearance of s	1	All spots					
cycle	Ω_0	В	Ω_{17}	Ω_0	В	Ω_{17}			
12	14.348 ± 0.016	3.004 ± 0.209	14.092 ± 0.003	14.337 ± 0.015	2.975 ± 0.200	14.081 ± 0.003			
13	14.430 ± 0.024	1.008 ± 0.277	14.344 ± 0.008	14.436 ± 0.015	1.147 ± 0.199	14.338 ± 0.002			
14	14.536 ± 0.014	2.724 ± 0.219	14.302 ± 0.005	14.606 ± 0.017	2.927 ± 0.275	14.353 ± 0.006			
15	14.485 ± 0.017	1.023 ± 0.209	14.398 ± 0.003	14.489 ± 0.022	1.277 ± 0.271	14.380 ± 0.003			
16	14.550 ± 0.011	2.630 ± 0.123	14.325 ± 0.003	14.549 ± 0.009	2.640 ± 0.109	14.323 ± 0.002			
17	14.288 ± 0.020	4.578 ± 0.214	13.899 ± 0.004	14.300 ± 0.011	4.680 ± 0.112	13.899 ± 0.003			
18	14.241 ± 0.017	3.924 ± 0.185	13.907 ± 0.003	14.229 ± 0.020	3.707 ± 0.201	13.912 ± 0.004			
19	14.217 ± 0.026	4.252 ± 0.231	13.849 ± 0.007	14.207 ± 0.017	4.044 ± 0.185	13.859 ± 0.003			
20	14.232 ± 0.027	2.646 ± 0.335	14.006 ± 0.005	14.254 ± 0.012	3.096 ± 0.141	13.989 ± 0.003			
21	14.236 ± 0.027	1.926 ± 0.293	14.071 ± 0.004	14.240 ± 0.018	2.010 ± 0.202	14.068 ± 0.003			
22	14.190 ± 0.027	1.420 ± 0.235	14.068 ± 0.007	14.158 ± 0.015	1.270 ± 0.121	14.055 ± 0.005			
23	14.460 ± 0.026	0.901 ± 0.309	14.380 ± 0.008	14.461 ± 0.024	0.912 ± 0.293	14.383 ± 0.003			

Table 1. Rotation parameters for individual cycles in units of deg/day.

Table 2. Rotation parameters for three-cycle intervals in units of deg/day.

	Northern hemisphere								
	Firs	t appearance of s	All spots						
cycle	Ω_0	В	Ω_{17}	Ω_0	В	Ω_{17}			
12-14	14.221 ± 0.023	4.014 ± 0.301	13.877 ± 0.007	14.206 ± 0.015	4.248 ± 0.237	13.846 ± 0.005			
15-17	14.381 ± 0.022	3.583 ± 0.297	14.078 ± 0.004	14.380 ± 0.019	3.516 ± 0.274	14.081 ± 0.004			
18-20	14.414 ± 0.015	3.516 ± 0.182	14.114 ± 0.001	14.426 ± 0.014	3.596 ± 0.168	14.117 ± 0.001			
21-23	14.579 ± 0.021	1.488 ± 0.260	14.452 ± 0.003	14.566 ± 0.019	1.322 ± 0.252	14.453 ± 0.003			
	Southern hemisphere								
	Firs	t appearance of s	_	All spots					
cycle	Ω_0	В	Ω_{17}	Ω_0	В	Ω_{17}			
12 - 14	14.372 ± 0.017	1.035 ± 0.240	14.284 ± 0.004	14.344 ± 0.020	0.604 ± 0.281	14.292 ± 0.004			
15-17	14.585 ± 0.020	3.168 ± 0.268	14.315 ± 0.003	14.571 ± 0.018	3.041 ± 0.243	14.311 ± 0.002			
18-20	14.513 ± 0.018	4.058 ± 0.252	14.168 ± 0.005	14.527 ± 0.009	4.517 ± 0.110	14.141 ± 0.001			
21-23	14.352 ± 0.007	3.068 ± 0.084	14.089 ± 0.001	14.337 ± 0.011	2.936 ± 0.135	14.087 ± 0.001			

Table 3. Rotation parameters for six-cycle intervals in units of deg/day.

	Northern hemisphere							
	Fir	st appearance of s	All spots					
cycle	Ω_0	В	Ω_{17}	Ω_0	В	Ω_{17}		
12-17	14.355 ± 0.012	3.116 ± 0.198	14.089 ± 0.004	14.305 ± 0.010	2.496 ± 0.171	14.091 ± 0.003		
18-23	14.588 ± 0.015	3.331 ± 0.196	14.303 ± 0.002	14.623 ± 0.009	2.687 ± 0.109	14.393 ± 0.001		
	Southern hemisphere							
	Fir	st appearance of s		All spots				
cycle	Ω_0	В	Ω_{17}	Ω_0	В	Ω_{17}		
12-17	14.611 ± 0.011	3.399 ± 0.141	14.3107 ± 0.002	14.614 ± 0.011	3.608 ± 0.161	14.305 ± 0.003		
18-23	14.347 ± 0.014	2.939 ± 0.193	14.095 ± 0.002	14.322 ± 0.011	2.806 ± 0.150	14.072 ± 0.001		

Table 4. Rotation parameters for 1878–1996 and 1874–2009 in units of deg/day.

	Northern hemisphere						
	Firs	t appearance of s	All spots				
time	Ω_0	В	Ω_{17}	Ω_0	В	Ω_{17}	
1878-1996	14.393 ± 0.008	3.309 ± 0.116	14.119 ± 0.003	14.414 ± 0.007	3.509 ± 0.119	14.114 ± 0.002	
1874-2009	14.424 ± 0.005	3.680 ± 0.072	14.111 ± 0.001	14.411 ± 0.003	3.455 ± 0.050	14.115 ± 0.001	
Southern hemisphere							
	Firs	t appearance of s		All spots			
time	Ω_0	В	Ω_{17}	Ω_0	В	Ω_{17}	
1878–1996	14.435 ± 0.008	4.931 ± 0.113	14.008 ± 0.002	14.413 ± 0.004	4.792 ± 0.067	14.003 ± 0.001	
1874-2009	14.420 ± 0.005	4.855 ± 0.073	14.006 ± 0.001	14.401 ± 0.004	4.652 ± 0.083	14.006 ± 0.001	



Fig. 2. Values of the merit function in color coding (in units of deg²) as a function of Ω_0 and *B* for all sunspots of the southern hemisphere in the limited region around the best-fit parameter values (**a**, *top*) for solar cycle 21; (**b**, *center*) for cycles 18–23; (**c**, *bottom*) for 1874–2009.

and 21, indicating the largest change in the Ω_0 value between two solar cycles.

The three-cycle Ω_0 values are mostly quite close to the average value of the corresponding three individual cycles. We note that this does not necessarily have to be so because the individual cycle fitting includes two more starting longitudes as additional free parameters. However, that the 1-cycle and three-cycle values agree with each other provides additional evidence of the robustness of the method and the long-term persistence of the active longitudes. Only one three-cycle average for cycles 18–20 deviates from the corresponding individual cycles. This could mean that a temporary phase shift may occur within this time. Interestingly, we also find (see Fig. 6) that the asymmetries for the individual cycles are lowest for these cycles.

Table 1 shows that the Ω_0 values for the firstly appeared sunspots are very close to those obtained for all sunspots. Except for the north of cycle 15 and for the south of cycles 14 and 22, they agree within the errors for all cycles and hemispheres. This close agreement is perhaps slightly surprising, but shows that the evolution of active longitudes is rather insensitive to how the sunspots are classified. Actually, in the database used, not all first appearances are new sunspots (because some were already borne on the invisible side) or firstly appeared sunspots (because some spots lived long enough to reappear). In addition, most sunspots do not live very long, experiencing only a small phase change due to their specific motion. These features alleviate any existing differences between the two sunspot classes. Table 5 shows that the parameters for all sunspots have slightly smaller average errors than the firstly appeared spots. This is particularly true for the basic fit parameters Ω_0 and *B*, but the errors of Ω_{17} approach each other for long time intervals.

Figure 4 depicts the optimal *B* parameters values using all sunspots for all individual cycles. The B values vary greatly from cycle to cycle, relatively much more than the Ω_0 parameters, attaining the lowest value of 0.467 during cycle 18 in the northern hemisphere and the largest value of 4.68 in cycle 17 in the south. The *B* values for firstly appeared sunspots agree within the error bars with those for all spots except for the north of cycle 15 where they differ by three standard deviations. This good agreement gives support that, despite their rather larger range, the optimal B values are reliable. We also note that for both hemispheres the *B* values calculated for the whole 136-year interval (see Table 4) are clearly above the mean *B* value for individual cycles. This is another demonstration that the optimal parameter values can, for the reasons discussed above, slightly vary depending on the time interval studied. Finally, Fig. 5 depicts the values of Ω_{17} calculated using Eq. (1) and the optimal values for Ω_0 and *B*. The panels are for the same cases as in Fig. 3 for Ω_0 .



Fig. 3. Values of Ω_0 using all sunspots for (*left panels*) individual cycles and (*right panels*) three-cycle intervals. The *upper (lower) panels* are for the northern (southern) hemisphere. The horizontal lines depict the best-fit value of Ω_0 for the whole interval.

Fig. 4. Values of the *B* parameter using all sunspots for (*left*) the northern and (*right*) southern hemisphere of all individual cycles. The horizontal lines depict the best-fit values of *B* for the whole interval.

(11)

Table 5. Average errors for different time intervals in units of deg/day.

E_{Ω_0}	bearance of E_B		nemisphere	All spots	
E_{Ω_0}		1	<i>г</i>	1	
0	E_B	Eo	Г		
0.021		-12 ₁₇	E_{Ω_0}	E_B	$E_{\Omega_{17}}$
±0.021	± 0.265	± 0.005	± 0.017	± 0.210	± 0.004
± 0.020	± 0.260	± 0.004	± 0.016	± 0.205	± 0.003
±0.014	± 0.197	± 0.003	± 0.010	± 0.140	± 0.002
± 0.009	± 0.124	± 0.002	± 0.007	± 0.119	± 0.002
± 0.005	± 0.072	± 0.001	± 0.003	± 0.050	± 0.001
Southern hemisphere					
First appearance of spots					
E_{Ω_0}	E_B	$E_{\Omega_{17}}$	E_{Ω_0}	E_B	$E_{\Omega_{17}}$
± 0.021	± 0.237	± 0.005	± 0.016	± 0.192	± 0.003
±0.016	± 0.200	± 0.003	± 0.015	± 0.192	± 0.002
±0.013	± 0.167	± 0.002	± 0.011	± 0.156	± 0.002
± 0.008	± 0.113	± 0.002	± 0.004	± 0.067	± 0.001
± 0.005	± 0.073	± 0.001	± 0.004	± 0.057	± 0.001
	$ \begin{array}{c} \pm 0.014 \\ \pm 0.009 \\ \pm 0.005 \end{array} \\ \hline First app \\ \hline E_{\Omega_0} \\ \pm 0.021 \\ \pm 0.016 \\ \pm 0.013 \\ \pm 0.008 \end{array} $	$\begin{array}{r} \pm 0.014 \\ \pm 0.197 \\ \pm 0.009 \\ \pm 0.124 \\ \pm 0.005 \\ \pm 0.072 \\ \hline \\ \hline \\ First appearance of \\ \hline \\ E_{\Omega_0} \\ \hline \\ \hline \\ \hline \\ E_{\Omega_0} \\ \hline \\ $	$\begin{array}{ccccccc} \pm 0.014 & \pm 0.197 & \pm 0.003 \\ \pm 0.009 & \pm 0.124 & \pm 0.002 \\ \pm 0.005 & \pm 0.072 & \pm 0.001 \\ \hline & & & \\ \hline E_{\Omega_0} & E_B & E_{\Omega_{17}} \\ \pm 0.021 & \pm 0.237 & \pm 0.005 \\ \pm 0.016 & \pm 0.200 & \pm 0.003 \\ \pm 0.013 & \pm 0.167 & \pm 0.002 \\ \pm 0.008 & \pm 0.113 & \pm 0.002 \\ \hline \end{array}$	$\begin{array}{c} \pm 0.014 \pm 0.197 \pm 0.003 \pm 0.010 \\ \pm 0.009 \pm 0.124 \pm 0.002 \pm 0.007 \\ \pm 0.005 \pm 0.072 \pm 0.001 \pm 0.003 \\ \hline \\ $	$ \begin{array}{c} \pm 0.014 \pm 0.197 \pm 0.003 \pm 0.010 \pm 0.140 \\ \pm 0.009 \pm 0.124 \pm 0.002 \pm 0.007 \pm 0.119 \\ \pm 0.005 \pm 0.072 \pm 0.001 \pm 0.003 \pm 0.050 \\ \hline \\ $

5. Non-axisymmetry

We define the measure of non-axisymmetry $\boldsymbol{\Gamma}$ in the same way as earlier

$$\Gamma = \frac{W_1 - W_2}{W_1 + W_2},\tag{9}$$

where W_1 and W_2 denote the normalized area of sunspots that appeared within (W_1) or outside (W_2) the two active longitudes, which are taken here as the two bands

$$W_1 = \sum_{k,i} W_{ik}, \text{ if } (|\lambda_{ik} - \Lambda_{ik}| < 45^\circ \text{ or } 360^\circ - |\lambda_{ik} - \Lambda_{ik}| < 45^\circ)$$
(10)

$$W_2 = \sum_{k,i} W_{ik}$$
, if $(|\lambda_{ik} - \Lambda_{ik}| > 45^\circ \text{ and } 360^\circ - |\lambda_{ik} - \Lambda_{ik}| > 45^\circ)$.

We have plotted the values of Γ for the individual cycles and the three-cycle periods in Fig. 6. One can see that there are large differences between the non-axisymmetries of the individual cycles. The largest non-axisymmetry for both hemispheres is found during cycle 16. For the northern hemisphere the nonaxisymmetry during this cycle is significantly larger than during any other cycle. We note that this cycle marked the beginning of the rise in solar activity that culminated with the Modern Grand Maximum of the last century. Non-axisymmetry is smallest in both hemispheres during cycle 19, the record cycle according to sunspot activity. This result agrees with Vernova et al. (2002), who noted that the excessive sunspot activity during cycle 19 was longitudinally symmetric, while some other cycles were stronger in longitudinally ordered sunspot activity.

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Fig. 5. Values of Ω_{17} using all sunspots for (left panels) individual cycles and (right panels) three-cycle intervals. The upper (lower) panels are for the northern (southern) hemisphere. The horizontal lines depict the best-fit value of Ω_{17} for the whole interval.

Fig. 6. Nonaxisymmetry parameter using all sunspots for (left panels) individual cycles and (right panels) three-cycle intervals. The upper (lower) panels are for the northern (southern) hemisphere.

While the non-axisymmetry does not depict any clear trend or oscillation pattern, the non-axisymmetries of the two hemispheres seem to roughly follow each other except during the last few years. This is also seen in the three-cycle averages of Fig. 6. Furthermore, in Table 6 we have given the average values of the non-axisymmetry parameters for the different time lengths. The non-axisymmetry slowly decreases as the time length increases. This decrease in non-axisymmetry with increasing time interval most likely results from the long-term variation in solar rotation. Table 6 shows that non-axisymmetries for all sunspots are slightly but consistently higher than for the first appearances of sunspots. This suggests that strong, long-living sunspots occur and remain preferentially around the two active longitudes.

6. Discussion

Let us first compare the present values of the rotation parameters with those obtained earlier for sunspots using the original, slightly less correct treatment of the differentially rotating coordinate system of active longitudes. Usoskin et al. (2005) used the first appearance of sunspot groups and found the best-fit values of $\Omega_0 = 14.33 \pm 0.01$ (deg/day) and $B = 3.40 \pm 0.03$ (deg/day) for the northern hemisphere and $\Omega_0 = 14.31 \pm 0.01$ (deg/day) and $B = 3.39 \pm 0.02$ (deg/day) for the southern hemisphere in 1878– 1996. The corresponding values obtained here (see Table 4) are $\Omega_0 = 14.393 \pm 0.008$ (deg/day) and $B = 3.309 \pm 0.116$ (deg/day) for the northern hemisphere, and $\Omega_0 = 14.435 \pm 0.008 (deg/day)$ and $B = 4.931 \pm 0.113$ (deg/day) for the southern hemisphere. Accordingly, the two values for Ω_0 and the value of *B* for the southern hemisphere obtained here are significantly larger than found earlier (Usoskin et al. 2005). The main difference between the two methods seems to be so that the best-fit solution now tends to emphasize the hemispheric difference, especially in the latitudinal variation parameter B, more strongly than earlier. This contradicts the result of Heristchi & Mouradian (2009), who found no north-south asymmetry in solar rotation. Large B values for both hemispheres strongly suggest that sunspot groups follow the differential rotation law, which differs from the conclusion of Balthasar (2007).

The present method leads to a larger asymmetry in the southern hemisphere than found earlier. The values of the Γ parameter obtained now (see Table 6) for the northern ($\Gamma = 0.110$) and southern hemisphere ($\Gamma = 0.109$) in 1878–1996 are roughly equal, while the asymmetry in the north ($\Gamma = 0.11$) obtained by Usoskin et al. (2005) was the same as here but larger than in the south ($\Gamma = 0.09$). This suggests that the rotation of the southern hemisphere may be more complicated or variable than in the northern hemisphere, and that the present improved method is indeed more accurate in treating such variable rotation of active longitudes. We note also that the non-axisymmetry parameters for the first appearances of sunspots in the last three solar cycles obtained here (0.17 for north and 0.13 for south, see Fig. 6) are significantly higher than those found by the earlier method (0.12 for north and 0.09 for south; Zhang et al. 2007b), giving further evidence for the superiority of the present method.

Table 6. Average value of non-axisymmetry Γ for different time intervals.

	Northern hemisph	ere	Southern hemisphere		
	First appearance of spots	All spots	First appearance of spots	All spots	
time	Γ	Г	Г	Г	
1 cycle	0.223	0.266	0.233	0.263	
3-cycle	0.153	0.163	0.155	0.188	
6-cycle	0.120	0.140	0.115	0.125	
1878–1996	0.110	0.112	0.109	0.119	
1874-2009	0.108	0.092	0.092	0.103	

In the present method, the values of Ω_0 are determined far more accurately than by the earlier method, but B is determined less accurately. Although no exact values for individual cycles were given by Usoskin et al. (2005), one can see in their Fig. 5 that the range of Ω_0 and B for the northern hemisphere of the individual cycles varied roughly between 13.65 and 14.7, and between 1.6 and 4.1, respectively, while here they were (see Table 1) roughly between 14.2 and 14.6, and between 0.4 and 4.3. However, despite these differences in the best-fit values of the two rotation parameters, the rotation rates at the maximum sunspot latitude are quite similar here ($\Omega_{17} = 14.119 \pm 0.003$ for north and $\Omega_{17} = 14.008 \pm 0.002$ for south) and earlier ($\Omega_{17} =$ 14.04 for north and $\Omega_{17} = 14.02$ for south). Thus we can verify the earlier observation that, over the whole interval, the mean rotation in the southern hemisphere is slower than in the north. The significance of this difference is considerably larger than found earlier. Pelt et al. (2007) also found that the average rotation was slower in the southern hemisphere.

As mentioned above, we used the improved method to study the differential rotation of active longitudes during the last three solar cycles using solar X-ray flares (Zhang et al. 2011). We found there that the new method increased the amount of nonaxisymmetry considerably and made the rotation parameters for three flare classes agree with each other. For instance, the for Cclass flares $\Omega_0 = 14.59 \pm 0.04$ and $B = 1.43 \pm 0.50$ for the northern hemisphere and $\Omega_0 = 14.39 \pm 0.04$ and $B = 2.80 \pm 0.50$ for the southern hemisphere. Comparing these with the present results depicted in Table 2, one finds excellent agreement for both parameters. The results for the two solar variables for the mean rotation rate at the latitude of maximum sunspot production also agree fairly closely, being $\Omega_{17} = 14.453 \pm 0.003/14.468 \pm 0.002$ for north and Ω_{17} = 14.087 \pm 0.001/14.151 \pm 0.002 for south for sunspots versus C-class flares. These results verify with exceedingly high statistical significance that the two solar hemispheres rotated, on an average, considerably differently during the last three solar cycles, the northern hemisphere much faster, the southern hemisphere slightly lower than the Carrington rotation rate.

Finally, we would like to note on the long-term evolution of the rotation parameters. As seen in Fig. 3, the Ω_0 values for the individual cycles depict relatively large and statistically highly significant temporal variations. Moreover, the long-term variations in Ω_0 are quite different in the two hemispheres. In the northern hemisphere, there is indication of a long-term increasing trend, upon which a multi-decadal oscillation is superimposed. In the southern hemisphere, there may be a small decreasing trend, which may be superimposed on a century-scale oscillation. Interestingly, the long-term variations in Ω_0 (both trends and oscillations) are roughly opposite in the northern and southern hemispheres. The *B* values also depict large long-term variability (see Fig. 4), but there does not seem to be any systematic long-term development in either hemisphere, either in trend or oscillation. However, as for Ω_0 , the *B* values depict roughly opposite long-term evolution in the two hemispheres, except during the two recent cycles.

We note that the long-term evolution of Ω_{17} (see Fig. 5) is quite similar to that of Ω_0 despite the rather unsystematic longterm change in the *B* parameter. This supports the view that the obtained results for the long-term evolution are indeed reliable and significant. Compared to Ω_0 , the long-term evolution of Ω_{17} for individual cycles seems to emphasize the increasing trend in the northern hemisphere, suppressing the amplitude of oscillation around the trend, while the centennial oscillation appears more systematically in the southern hemisphere in Ω_{17} than in Ω_0 . The three-cycle values also underline the opposite temporal evolutions in the two hemispheres with increasing rotation rate in the north and decreasing rotation in the south.

Earlier studies have already noted that the rotation parameters depict long-term trends and long-term oscillations (see, e.g., Makarov et al. 1997; Kitchatinov et al. 1999; Javaraiah et al. 2005; Brajsa et al. 2006, and references therein). For instance, Javaraiah et al. (2005) found that the Ω_0 parameter depicts a fairly systematic decline during cycles 12 to 23 from roughly $\Omega_0 = 14.55 \text{ deg/day}$ to $\Omega_0 = 14.45 \text{ deg/day}$, with some evidence of weak century-scale oscillation superimposed. The *B* parameter was also found to have a small decreasing trend from about B = 2.7 to about B = 2.4, but the evolution of *B* was dominated by a roughly 80-year oscillation with an amplitude of $\pm 0.3 \text{ deg/day}$. However, these results cannot be directly compared with the long-term evolution of the best-fit values of the two rotation parameters found in this study since Javaraiah et al. (2005) did not separately consider the two hemispheres.

As noted above, the largest non-axisymmetry for both hemispheres is found for cycle 16, which started the rise of solar activity in the 20th century. On the other hand, the smallest asymmetry is found for cycle 19, which was the most active cycle, starting the decrease in sunspot activity thereafter. Although the values of Ω_0 or Ω_{17} were considerably larger for cycle 16 than cycle 19 in both hemispheres, they were not exceptional or extremum values for either of these two cycles. Accordingly, the relation between solar rotation and solar activity seems to be rather complicated.

Berdyugina et al. (2006) and Usoskin et al. (2007) summarized the dynamo modes related to the occurrence of preferred longitudes in solar and stellar activity and discussed the interpretation of the existence of active longitudes within several solar dynamo related and other scenarios (see also Ruzmaikin 1998; Rädler et al. 1990; Moss 1999, 2004; Mason et al. 2002; Fluri & Berdyugina 2005; Bassom et al. 2005; Korhonen & Elstner 2005). Berdyugina et al. (2006) noted that in most dynamo models the generated structures are expected to rotate quasirigidly, while the migration of active longitudes could be caused by a stroboscopic effect. Nevertheless, the significantly varying rate of asymmetry and rotation from one cycle to another, as well as the significant hemispheric differences provide observational constraints that need to be taken into account in dynamo modeling.

7. Conclusions

Using a refined version of the dynamic, differentially rotating coordinate system, we have determined the optimal values of the two rotation parameters Ω_0 and *B* (see Eq. (1)) and the average rotation rate Ω_{17} for all sunspot groups and the first appearances of sunspots in the two hemispheres separately. We have also calculated these parameters for different lengths of time. We list below our main results and findings:

- 1. We found the following parameter values for, e.g., all sunspots over the whole 136-year time interval: $\Omega_0 = 14.411 \pm 0.003$ (deg/day) and $B = 3.455 \pm 0.050$ (deg/day) for the northern hemisphere and $\Omega_0 = 14.401 \pm 0.004$ (deg/day) and $B = 4.652 \pm 0.083$ (deg/day) for the southern hemisphere. Typically, the refined method yielded somewhat larger values for Ω_0 than the earlier method (Usoskin et al. 2005).
- 2. The values of Ω_0 were determined much more accurately but *B* slightly less accurately than earlier (Usoskin et al. 2005). However, the average rotation rates Ω_{17} at the maximum sunspot latitude of 17° in 1878–1996 were found to be almost similar to earlier values.
- 3. The refined method yielded closely similar asymmetries in the two hemispheres, and a larger asymmetry in the southern hemisphere in 1878–1996 than earlier (Usoskin et al. 2005). In addition, the non-axisymmetry of the first appearances of sunspots in the last three solar cycles are significantly higher here than those found by the earlier method (Zhang et al. 2007b). This gives further evidence of the credibility and importance of the improvement of the method.
- 4. The optimal values of the two differential rotation parameters obtained here for all sunspots are in close agreement with those for first appearance spots in both hemispheres and all time intervals.
- 5. The two rotation parameters for sunspots during the last three solar cycles agree closely with the parameters obtained for different classes of X-ray flares using the same refined method. To our knowledge, such good agreement between rotation parameters of sunspots and flares in both hemispheres is unique.
- 6. Non-axisymmetry in sunspots during the last three solar cycles was somewhat smaller than in X-ray flares. Also, non-axisymmetry in all sunspots was systematically larger than for the first appearances of spots. These results support the view that strong and persistent forms of solar activity follow more closely to the two active longitudes.
- 7. The largest non-axisymmetry for both hemispheres is found during cycle 16. It is interesting to note that this cycle started the rise of solar activity to the Modern Grand Maximum of the last century. Non-axisymmetry is smallest in both hemispheres during cycle 19, the record cycle according to sunspot activity.
- 8. We found that the errors in optimal parameters decreased slowly but systematically with increasing time interval, implying that there is a fairly good temporal consistency in the rotation parameters.
- 9. The refined method emphasizes the hemispheric differences in the two parameters. Over the whole interval, the mean rotation in the southern hemisphere is slower than in the north. The significance of this difference is considerably larger than earlier. There is overwhelming agreement from sunspots and flares that the northern hemisphere rotated much faster than and the southern hemisphere slightly slower than the Carrington rotation rate during the last three solar cycles.

10. There is significant temporal variability in the two rotation parameters over the 136-year interval. In the northern hemisphere, Ω_0 and Ω_{17} follow long-term increasing trends, upon which a weak multi-decadal oscillation may be superimposed. In the southern hemisphere, Ω_0 and Ω_{17} exhibit decreasing trends, superimposed on a significant century-scale oscillation. Interestingly, the long-term variations in Ω_0 and Ω_{17} (both trends and oscillations) are roughly opposite in the northern and southern hemispheres.

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